



**SORET AND DUFOUR EFFECTS ON CONVECTIVE HEAT AND MASS TRANSFER FLOW OF A VISCOUS FLUID IN THE PRESENCE OF THERMOPHORESIS PARTICLE DEPOSITION**

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**ABSTRACT**

*We analyze the combined influence of thermal radiation, Soret and Dufour effect on convective heat and mass transfer flow of a viscous incompressible fluid in the presence of suction\injection and thermophoresis deposition particle. The non-linear equations governing the momentum, temperature and concentration along with the boundary conditions are solved by using finite element method with Mathematica6.1 software.*

**Keywords:** *Soret and Dufour effects, Thermophoresis, Non-Uniform Heat Source.*

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**1. INTRODUCTION**

Thermophoresis phenomenon has many practical applications in removing small particles from gas streams, in determining exhaust gas particle trajectories from combustion devices, and in studying the particulate material deposition on turbine blades. It has been found that thermophoresis is the dominant mass transfer mechanism in the modified chemical vapor deposition process as currently used in the fabrication of optical fiber performs. Thermophoretic deposition on radioactive particles is considered to be one of the important factors causing accidents in nuclear reactors. Keeping the applications in view several researchers (Goren[12], Talbot *et al.* [25], Homsy *et al.* [14], Epstein *et al.* [10], Garg and Jayataj[11], Jia *et al.*[13] Chiou[8], Selim *et al.* [24], Chamkha *et al.* [5] Chamkha and Pop [6], Seddeek [23], Chamkha *et al.* [7], Partha [20], Partha [21], Beg *et al.* [4], PuviArasuet *et al.* [22], Kandasamy *et al.* [16], Alam *et al.* [1], Alam *et al.* [2] Dulal Pal *et al.* [9]) have studied the effect of thermophoresis on convective heat and mass transfer flow in different configurations.

The effect of radiation on MHD flow and heat transfer problem has become more important industrially. At high temperature, radiation effect can be quite significant. Many processes in engineering areas occur at high temperature and knowledge of radiation heat transfer becomes very important for the design of the pertinent equipment. Nuclear power plants, gas turbines and the various propulsion devices for aircraft, missiles, satellites and space vehicles are examples of such engineering areas. Several authors (Bestman [3], Mankinde [19], Ibrahim *et al.* [15], Madhusudana *et al.* [18]) have investigated the effect of thermal radiation on heat and mass transfer flow under varied conditions.

In this paper we analyze the combined influence of thermal radiation, Soret and Dufour effect on convective heat and mass transfer flow of a viscous incompressible fluid in the presence of suction/injection and thermophoresis deposition particle. This problem has many applications, such as, technological and manufacturing engineering, radioactive particle deposition in nuclear reactors, deposition of silicon thin films and polymer extrusion process. The non-linear equations governing the momentum, temperature and concentration along with the boundary conditions are solved by using finite element method with Mathematica6.1 software.

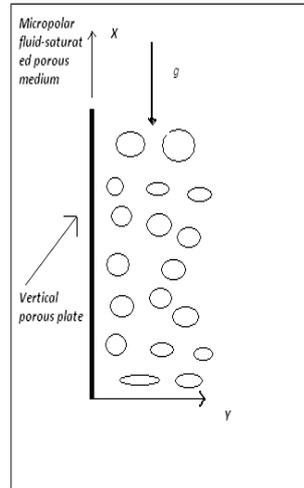
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## 2. MATHEMATICAL FORMULATION

We consider the two-dimensional study mixed convective heat and mass transfer flow of a viscous incompressible, non-conducting, fluid through a porous medium in the presence of thermophoresis as shown in fig.1. The x-axis is directed along the vertical surface and the y-axis is perpendicular to this as shown in figure 1. The vertical surface and the fluid are maintained same temperature and concentration initially. Instantaneously they raised to a temperature  $T_w (> T_\infty)$  and concentration  $C_w (> C_\infty)$  which remain unchanged. The effects of thermophoresis are being taken in the diffusion equation to help in the understanding of the mass deposition variation on the surface. The temperature gradient in the y-direction is much larger than that in x-direction and hence thermophoretic velocity component normal to the surface is of more importance.



**Figure-1:** Physical Model of the Coordinate System

Under the above stated physical situation, the governing boundary-layer and Darcy-Boussinesq approximations, the basic equations are given by:

$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} = 0 \quad (1)$$

$$u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} = \nu \frac{\partial^2 u}{\partial y^2} + k_1 \frac{\partial w}{\partial y} + \beta g (T - T_\infty) + \beta^* g (C - C_\infty) - \frac{\nu}{k} u - \frac{\sigma \mu_e^2 H_o^2}{\rho_o} u \quad (2)$$

$$u \frac{\partial w}{\partial x} + v \frac{\partial w}{\partial y} = \frac{k}{\rho j} \frac{\partial^2 w}{\partial y^2} - \frac{k}{\rho j} \left( 2w + \frac{\partial u}{\partial y} \right) \quad (3)$$

$$u \frac{\partial T}{\partial x} + v \frac{\partial T}{\partial y} = \frac{k_f}{C_p} \frac{\partial^2 T}{\partial y^2} + \frac{\nu}{C_p} \left( \frac{\partial u}{\partial y} \right)^2 + \frac{1}{C_p} q'' + \frac{D_m K_T}{C_s C_p} \frac{\partial^2 C}{\partial y^2} - \frac{1}{C_p} \frac{\partial (q_R)}{\partial y} \quad (4)$$

$$u \frac{\partial C}{\partial x} + v \frac{\partial C}{\partial y} = D_m \frac{\partial^2 C}{\partial y^2} + \frac{D_m K_T}{T_m} \frac{\partial^2 T}{\partial y^2} - \frac{\partial (V_T C)}{\partial y} \quad (5)$$

The associated boundary conditions on the vertical surface are defined as follows,

$$u = U_w = ax + L \frac{\partial u}{\partial y}, v = v_w(x), T = T_w, C = C_w \quad \text{at } y = 0 \quad (6)$$

$$u \rightarrow 0, v \rightarrow 0, T \rightarrow T_\infty, C \rightarrow C_\infty \quad \text{as } y \rightarrow \infty$$

In the above equations x and y represents coordinate axes along the continuous surface in the direction of motion and perpendicular to it, u and v are the velocity components along x and y directions, respectively. We note that  $V_0 < 0$  represents suction while  $V_0 > 0$  represents injection into the flow regime. Here we confine our attention to the suction ( $V_0 > 0$ ) of fluid through the porous medium.

The effect of thermophoresis is usually prescribed by means of an average velocity acquired by small particles to the gas velocity when exposed to a temperature gradient. In boundary layer flow, the temperature gradient in y-direction is very much larger than in the x-direction and therefore only the thermophoretic velocity in y-direction is considered. As a consequence, the thermophoretic velocity  $V_T$ , which appears in equation (4), is expressed as

$$V_T = -\frac{k_1 \nu}{T_r} \frac{\partial T}{\partial y} \quad (7)$$

in which  $k_1$  is the thermophoretic coefficient and  $T_r$  is the reference temperature. A thermophoretic parameter  $\tau$  is given by the relation

$$\tau = -\frac{k_f (T_w - T_\infty)}{T_r} \quad (8)$$

Where the typical values of  $\tau$  are 0.01, 0.1 and 1.0 corresponding to approximate Values of  $-k_1 (T_w - T)$  equal to 3, 30, 300K for a reference temperature of  $T=300K$ .

The non-dimensional heat/sink,  $q'''$  (19) is modeled as

$$q''' = \frac{k_f u_w(x)}{\nu x} (A1(T_w - T_\infty)u + B1(T - T_\infty)) \quad (9)$$

Where  $A1$  and  $B1$  are the coefficient of space and temperature dependent heat source/sink respectively. Here we make a note that the case  $A1 > 0$ ,  $B1 > 0$  corresponds to internal heat generation and that  $A1 < 0$ ,  $B1 < 0$  corresponds to internal heat absorption.

The energy equation with thermal radiation reduces in to

$$\rho C_p \left( u \frac{\partial T}{\partial x} + w \frac{\partial T}{\partial z} \right) = k_f \frac{\partial^2 T}{\partial y^2} + \mu \left( \frac{\partial u}{\partial y} \right)^2 + \frac{k_f U_0}{\nu x} (A1(T_w - T_\infty)u + B1(T - T_\infty)) + \frac{D_m K_T}{C_s C_p} \frac{\partial^2 C}{\partial y^2} + \frac{16\sigma^* T_\infty^3}{3\beta_R} \frac{\partial^2 T}{\partial y^2} \quad (10)$$

The mathematical analysis of the problem is simplified by introducing the following dimensionless functions  $f$ ,  $h$ ,  $\theta$ ,  $\phi$  and the similarity variable  $\eta$

$$\eta = \sqrt{\frac{U_w}{2\nu x}} y, \quad u = U_w f'(\eta), \quad v = -\left(\frac{\nu U_w}{2x}\right)^{1/2} (f(\eta) + \eta f'(\eta)), \quad \theta = \frac{T - T_\infty}{T_w - T_\infty}, \quad \phi = \frac{C - C_\infty}{C_w - C_\infty}, \quad W = -\sqrt{\frac{U_w}{2\nu x}} U_w h(\eta) \quad (11)$$

Substituting the equation (11) in equations (2), (3), (5) and (11)

We obtain

$$f''' + B_1 h' + ff'' + G_r (\theta + N\phi) - \left(\frac{D^{-1}}{R_e} + M^2\right) f' = 0 \quad (12)$$

$$\lambda h'' - 2 \frac{\lambda}{G_1} (2h + f'') + (fh' + f'h) = 0 \quad (13)$$

$$\theta'' + Pr f \theta' + Pr Du \phi'' + \frac{4Rd}{3} \theta'' + Pr (A_1 f' + B_{11} \theta) + Pr Ec (f'')^2 = 0 \quad (14)$$

$$\phi'' + Sc f \phi' + Sc Sr \theta'' - \tau (\theta' \phi' + \theta'' \phi) = 0 \quad (15)$$

Where

$$G_r = \frac{2\beta g x (T_w - T_\infty)}{U_w^2} (\text{Local Grashof Number}), \quad N = \frac{\beta^* (C_w - C)}{\beta (T_w - T)} (\text{Buoyancy ratio}),$$

$$\begin{aligned}
 D^{-1} &= \frac{U_w k_p}{2\nu x} \text{ (Inverse Darcy parameter),} & M^2 &= \frac{2\sigma\mu_e^2 H_o^2 x}{\rho U_w} \text{ (Hartmann number)} \\
 Re &= \frac{U_w L^2}{2\nu x} \text{ (Local Reynolds number),} & Pr &= \frac{\mu C_p}{k_f} \text{ (Pr andtl Number),} \\
 Sc &= \frac{\nu}{D_m} \text{ (Schmidt number),} & Du &= \frac{D_m K_T (C_w - C_\infty)}{C_s C_p (T_w - T_\infty)} \text{ (Dufour parameter),} \\
 Sr &= \frac{D_m (T_w - T_\infty)}{\nu T_m (C_w - C_\infty)} \text{ (Soret parameter).} & Ec &= \frac{U_w^2}{C_p (T_w - T_\infty)} \text{ (Ec ker t Number),} \\
 Rd &= \frac{4\sigma^* T^3}{\beta_R k_f} \text{ (Radiation parameter) ,} & B1 &= \frac{k_1}{\nu} \text{ (Coupling constant parameter)} \\
 \lambda &= \frac{\gamma}{\rho \nu j} \text{ (Micro polar parameter),} & G1 &= \frac{\gamma}{\nu k} \text{ (Micro polar parameter)} \\
 \nu &= \frac{\mu + k}{\rho} \text{ (Kinematic viscosity),} & k_1 &= \frac{k}{\rho} (k_1 > 0) \text{ (Coupling constant)} \\
 u &= -\psi_y & v &= \psi_x .
 \end{aligned}$$

The corresponding transformed boundary conditions are

$$\begin{aligned}
 f'(0) = 1 + Af'', f(0) = V_0 \sqrt{\frac{2x}{\nu U_w}} = fw, \theta(0) = 1, \phi(0) = 1 \text{ at } \eta = 1 \\
 f'(\infty) = 0, \theta(\infty) = 0, \phi(\infty) = 0 \text{ as } \eta \rightarrow \infty
 \end{aligned} \tag{16}$$

### 3. METHOD OF SOLUTION

The non linear coupled equations governing the flow have been solved by using FEM with quadratic approximation polynomials. The resulting local stiffness matrices are assembled by using inter element continuity, equilibrium and boundary conditions.

### 4. SHEAR STRESS, NUSSELT NUMBER AND SHERWOOD NUMBER

The major physical quantities of interest in this problem are the local skin friction coefficient  $C_f$ , the local Nusselt number  $Nu_x$  and the local Sherwood number  $Sh_x$  are defined, respectively by

$$C_f = \frac{f''(0)}{Re_x^{1/2}}, \quad Nu_x = -\frac{\theta'(0)}{Re_x^{1/2}}, \quad Sh_x = -\frac{\phi(0)}{Re_x^{1/2}}$$

The correctness of the current numerical method is checked with the results obtained by Loganathan *et al.* [2010] for the local skin friction coefficient  $C_f$ , local Nusselt number  $Nu_x$  and local Sherwood number  $Sh_x$  for various values of thermophoretic parameter  $\tau$  and suction parameter  $V_0$ . Thus, it is seen from Table 1 that the numerical results are in close agreement with those published previously.

### COMPARISION

Comparison of results of  $C_f, Nu_x, Sh_x$  with previously published data,  $Du = 0, Sr = 0, Ec = 0, A1 = 0, B11 = 0, Rd = 0$

**Table-1**

Du	Sr	$Nu_x$	$Sh_x$
0.5	0.5	-0.952739	-0.514925
	1.0	-0.981055	-0.439975
	1.5	-1.01065	-0.359391
	2.0	-1.04169	-0.2727
1	0.5	-0.870524	-0.521268
	1.0	-0.915672	-0.452475
	1.5	-0.966592	-0.373596
	2.0	-1.02494	-0.282221
1.5	0.5	-0.786301	-0.527484
	1.0	-0.844034	-0.465539
	1.5	-0.914626	-0.389352
	2.0	-1.00453	-0.292617
2	0.5	-0.700188	-0.53357
	1.0	-0.765569	-0.479223
	1.5	-0.852782	-0.407087
	2.0	-0.979596	-0.304151
2.5	0.5	-0.612292	-0.539526
	1.0	-0.679561	-0.49361
	1.5	-0.778022	-0.427486
	2.0	-0.94833	-0.317407

**Table-2**

Parameter		Loganathan		
		$C_f$	$Nu_x$	$Sh_x$
$\tau$	1.0	6.034338	-2.580389	-5.260410
	2.0	6.034184	-2.580390	-5.804620
	3.0	6.034067	-2.580391	-8.366696
$V_0$	0.5			
	1.0			
	1.5			
	4.0	7.899134	-7.543495	-4.788675
	5.0	9.665340	-9.336595	-5.882326
	6.0	11.45915	-11.14177	-6.988321
Parameter		Present results		
		$C_f$	$Nu_x$	$Sh_x$
$\tau$	1.0	6.034336	-2.580385	-5.260409
	2.0	6.034189	-2.580385	-5.804615
	3.0	6.034069	-2.580383	-8.366989
$V_0$	0.5	4.125482	-5.026587	-1.236587
	1.0	5.354135	-6.154873	-2.365989
	1.5	6.054877	-6.998514	-3.654871
	4.0	7.899139	-7.543493	-4.788673
	5.0	9.665336	-9.336592	-5.882322
	6.0	11.45919	-11.14169	-6.988319

Computed values of local Nusselt number and local Sherwood number for

$$V_0 = 0.5; P_r = .71; S_c = 0.22; G_r = 5; G_m = 5; D_a = 0.1; r = 10; E_c = 0.01; \lambda = 0.1; s = 0.1; G1 = 0.5; B1 = 0.5; \tau = 0.1.$$

## 5. DISCUSSION OF THE RESULTS

In this analysis we discuss the effect of thermal radiation, dissipation, Soret and Dufour on convective heat and mass transfer flow of a viscous, electrically conduction fluid past a stretching sheet in the presence of non-uniform heat source. The results are presented graphically in figures.2a-7a for different parametric variations. Comparison of the present results with previous works is performed and excellent agreements have been obtained. The non-linear coupled differential equations are solved by Galerkin finite analysis with three noded line segments. In the absence of the results are compared with Loganathan [17].

Figs.2a-2d shows the influence of dissipation on the velocity, Angular velocity, temperature and concentration. It is pointed out that the presence of Eckert number increases the velocity, Angular velocity and temperature. This is due to the fact that the thermal energy is reserved in the fluid on account of friction heating. Hence the velocity and temperature rises in the entire boundary layer. However, the mass concentration is to reduce marginally with increase in  $Ec$  (fig.2c).

Figs.3a-3d and 4a-4d represents the variation of velocity, Angular velocity, temperature and concentration with space dependent heat source/temperature dependent heat source in the boundary layer. An increase in the space dependent /temperature dependent heat source increases the velocity, Angular velocity and temperature in the boundary layer. This may be attributed to the fact that the presence of the heat source generates energy in the boundary layer and as a consequence the velocity and temperature rises. In the case of heat absorption  $A1 < 0$ ,  $B11 < 0$ , the velocity and temperature falls, owing to the absorption of energy in the boundary layer. In the case on concentration we find that it reduces with  $A1 > 0$ ,  $B11 > 0$  and increases with  $A1 < 0$ ,  $B11 < 0$  in the boundary layer.

Figs.5a-5d depicts the influence of Radiation parameter on the velocity, Angular velocity, temperature and concentration. It is observed that there is a significance rise in the velocity in the presence of thermal radiation throughout the boundary layer. Further increase in the values of thermal radiation parameter results in the increase of the boundary layer. The presence of the thermal radiation is very significant on the variation of temperature. It is seen that the temperature rapidly increases in the presence of thermal radiation parameter throughout the thermal boundary layer. This may be attributed to the fact as the Rosseland radiative absorption parameter  $R^*$  diminishes the corresponding heat flux diverges and thus rises the rate of radiative heat transfer to the fluid causing a rise in the temperature of the fluid. The thickness of the boundary layer also increases in the presence of  $Rd$ . The effect of  $Rd$  on concentration is to diminish it in the solutal boundary layer.

The effect of Soret and Dufour effects on the velocity, Angular velocity, temperature and concentration is exhibited in figs.6a-6d. From fig.6a and 6b we find that the velocity and Angular velocity profiles experience an enhancement with increasing values of  $Sr$  (or decreasing Dufour parameter ( $Du$ )). This is due to the fact that an increase in  $Sr$  (or decrease in  $Du$ ) increases the thickness of the boundary layer. From fig.6c we find that an increase in  $Sr \leq 1.0$  (or  $Du \leq 0.04$ ) increases the temperature and for higher  $Sr \geq 1.5$  (or  $Du \geq 0.05$ ) we notice a decrease in temperature in the thermal boundary layer. An increase in  $Sr$  (or decrease in  $Du$ ) enhances the concentration. This may be attributed to the fact that the thickness of the solutal boundary layer increases with increase in  $Sr$  which in turn rises the concentration in the flow region (fig.6d).

Figs.7a-7d represents the influence of thermophoresis on velocity, Angular velocity, temperature and concentration. It is observed from fig.7a that increasing the values of thermophoresis parameter  $\tau$  decreases the thickness of the momentum boundary layer while an increase in  $\tau$ , increases the thickness of the thermal and solutal boundary layers (figs.7c-7d). From figure 7b we notice a reduction in angular velocity with increase in  $\tau$ .

The skin friction ( $\tau_1$ ) at the wall  $\eta=0$  is shown in tables 3 and 4 for different values of the parameters,  $A1$ ,  $B11$ ,  $Rd$ ,  $Sr$  &  $Du$ ,  $Ec$ . Here the thermo-diffusion smaller the skin friction. Increasing Soret parameter  $Sr$  (or decreasing Dufour parameter  $Du$ ) smaller the skin friction. An increasing in the heat generating source enhances  $\tau$  and reduces in the case heat absorbing source. An increase in micro polar parameter  $G1$ ,  $B1$  and  $\lambda$  leads to a reduction in  $\tau_1$ .

The rate of heat transfer (Nusselt number) at the wall  $\eta=0$  is shown in tables 3 and 4 evaluated for different parametric variation. Increasing  $Sr$  (or decreasing  $Du$ ) leads to an enhancement  $Nu$  at the wall. Higher the dissipation/Radiative heat flux lesser the Nusselt number at the wall. With reference to the heat source parameter, we find that the rate of heat transfer reduces with increase in heat generating source and enhances with heat absorption source.  $Nu$  increases with  $G1$  and decreases with  $B1$  and  $\lambda$ .

The rate of mass transfer (Sherwood number) at the wall  $\eta=0$  is shown in tables 3 and 4 evaluated with variations in different parameters. It can be seen from the tabular values that the rate of mass transfer increases with increase in  $Ec$ . Higher the radiative heat flux larger the rate of mass transfer at the wall. We noticed an enhancement in  $Sh$  increasing  $Sr$  (or  $Du$ ) smaller the rate of mass transfer at the wall. The rate of mass transfer decreases with space dependent heat source and reduces with temperature dependent heat source.  $Sh$  experiences an enhancement with  $G1$ ,  $B1$  and  $\lambda$ .

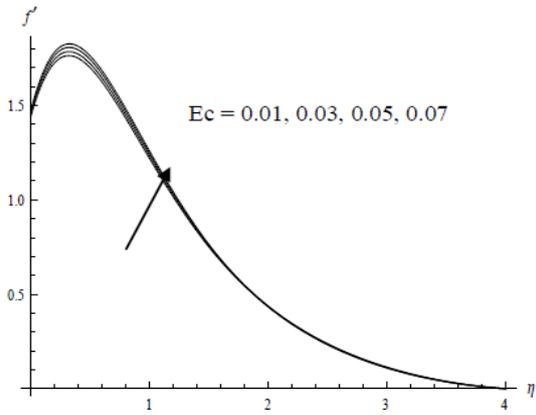


Figure-2a: Variation of  $f'$  with  $Ec$

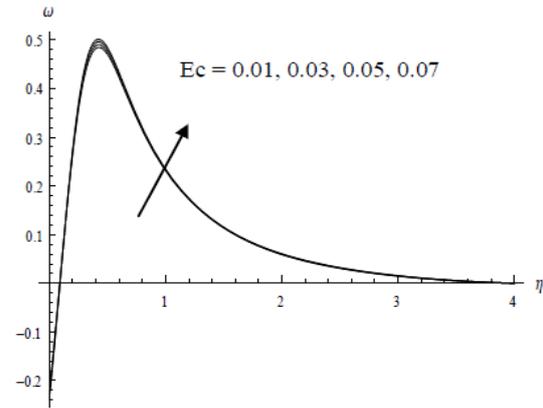


Figure-2b: Variation of  $\omega$  with  $Ec$

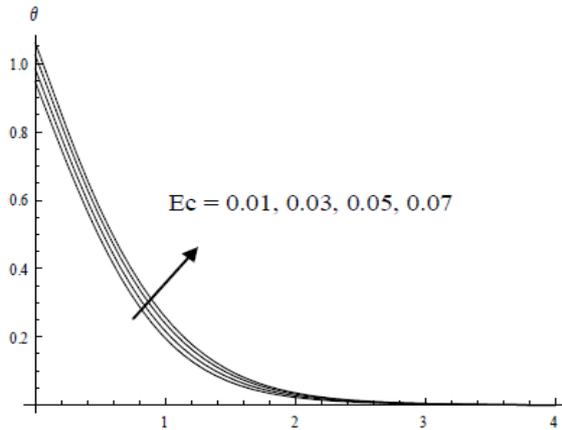


Figure-2c: Variation of  $\theta$  with  $Ec$

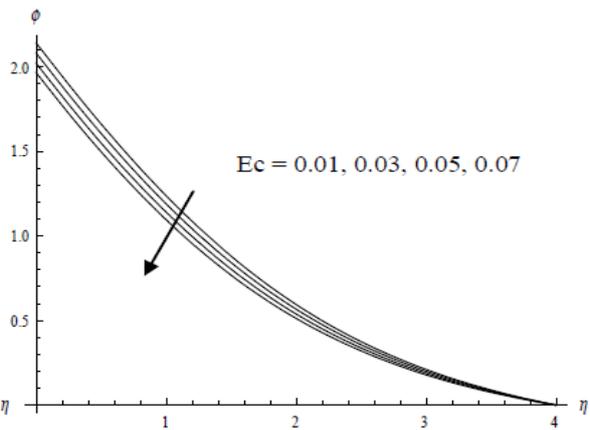


Figure-2d: Variation of  $\phi$  with  $Ec$

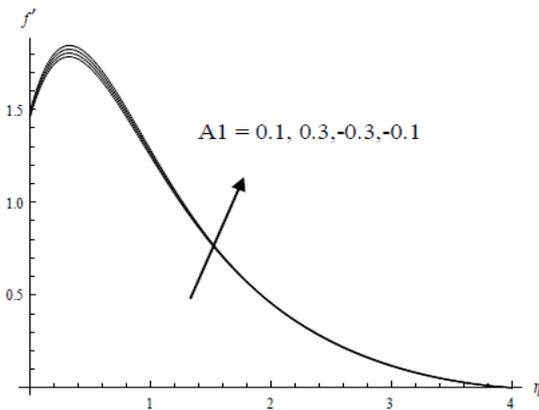


Figure-3a: Variation of  $\omega$  with  $A1$

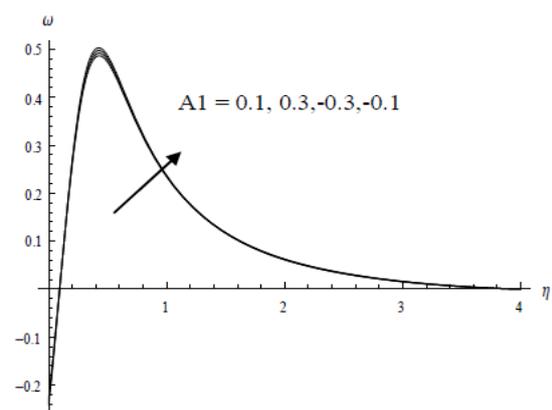


Figure-3b: Variation of  $\omega$  with  $A1$

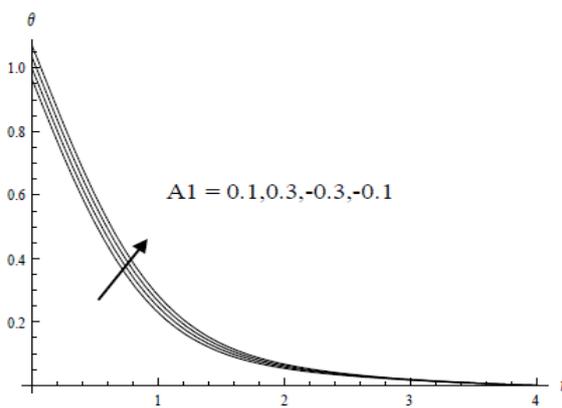


Figure-3c: Variation of  $\theta$  with  $A1$

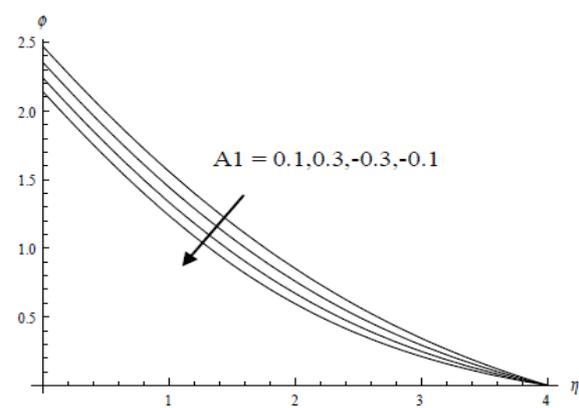


Figure-3d: Variation of  $\phi$  with  $A1$

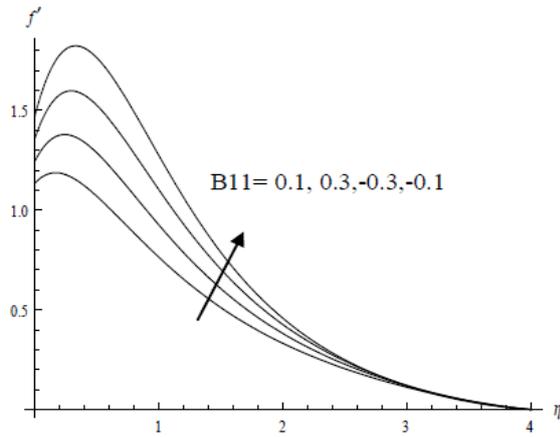


Figure-4a: Variation of  $f'$  with  $B_{11}$

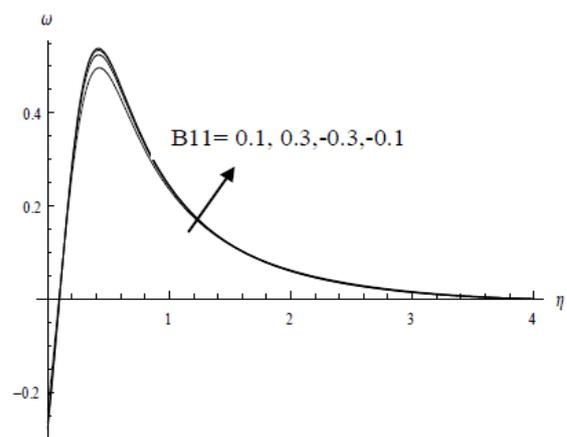


Figure-4b: Variation of  $\omega$  with  $B_{11}$

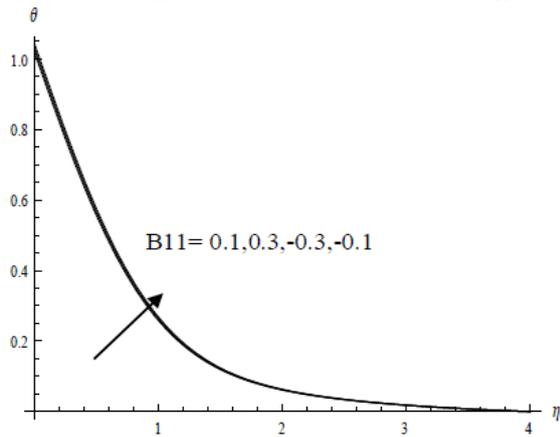


Figure-4c: Variation of  $\theta$  with  $B_{11}$

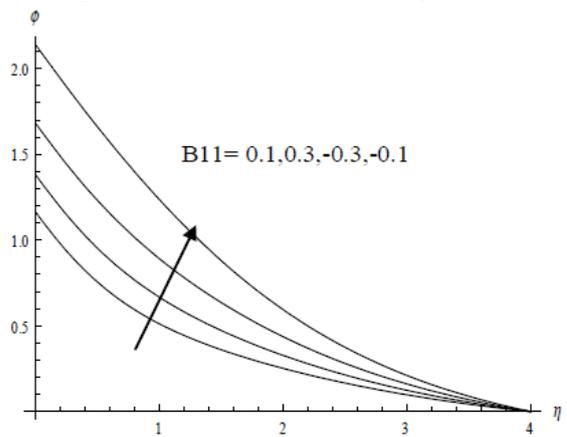


Figure-4d: Variation of  $\phi$  with  $B_{11}$

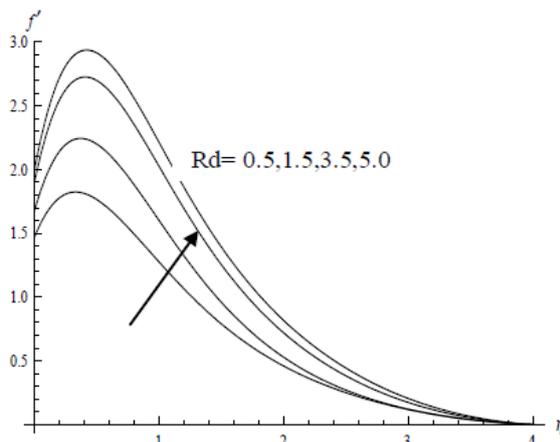


Figure-5a: Variation of  $f'$  with  $Rd$

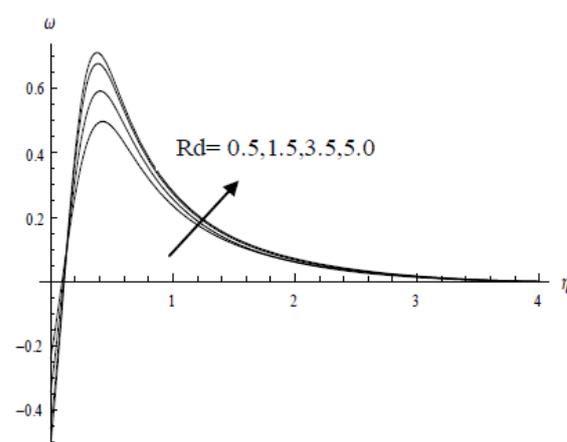


Figure-5b: Variation of  $\omega$  with  $Rd$

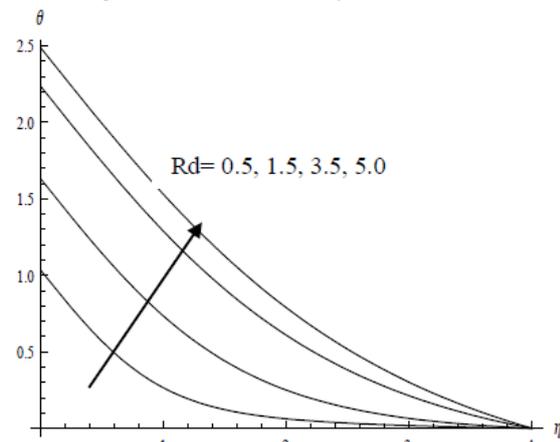


Figure-5c: Variation of  $\theta$  with  $Rd$

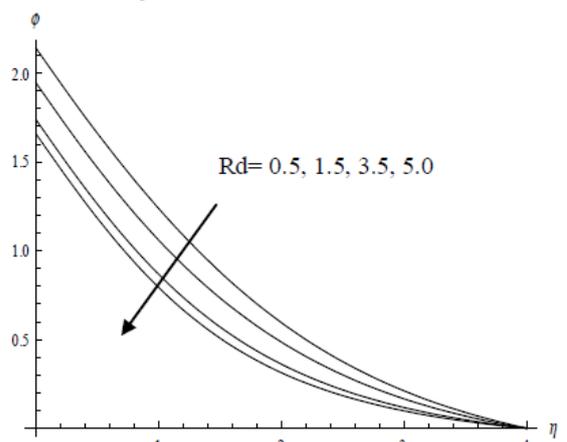


Figure-5d: Variation of  $\phi$  with  $Rd$

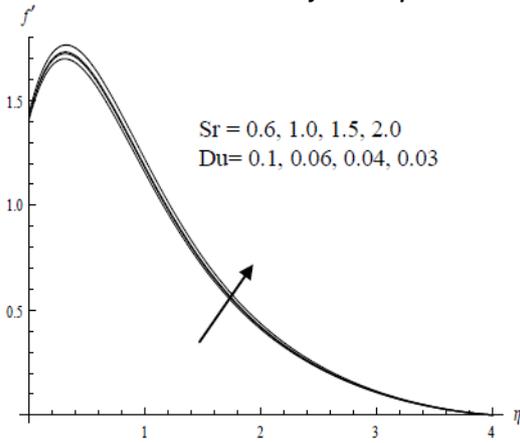


Figure-6a: Variation of  $f'$  with Sr & Du

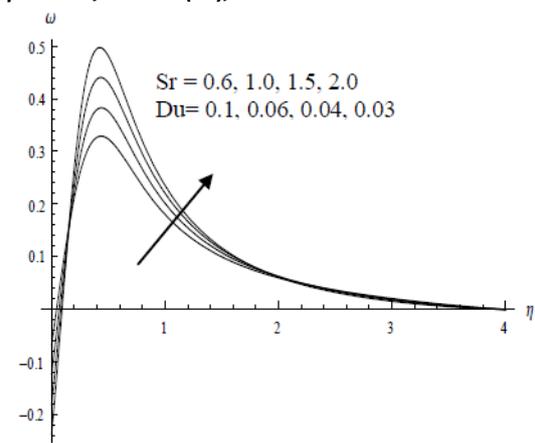


Figure-6b: Variation of  $\omega$  with Sr & Du

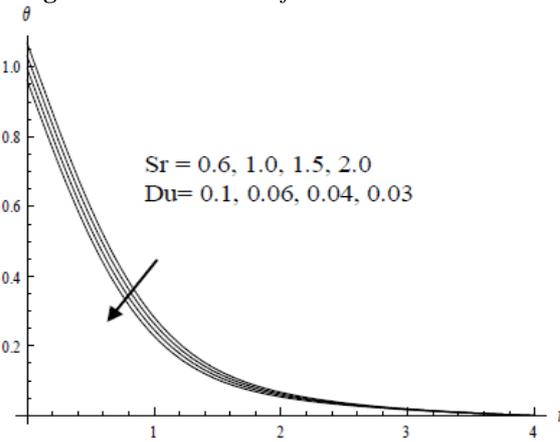


Figure-6c: Variation of  $\theta$  with Sr & Du

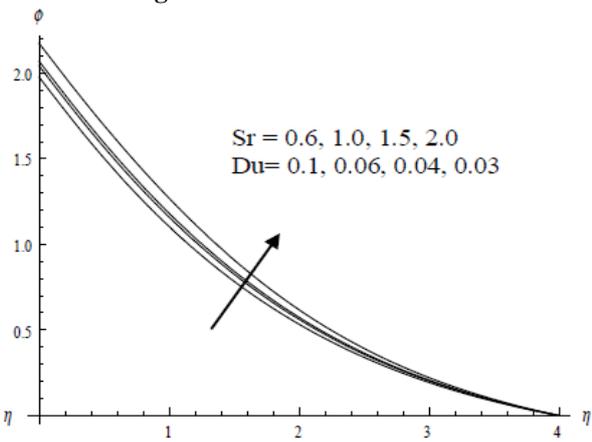


Figure-6d: Variation of  $\phi$  with Sr & Du

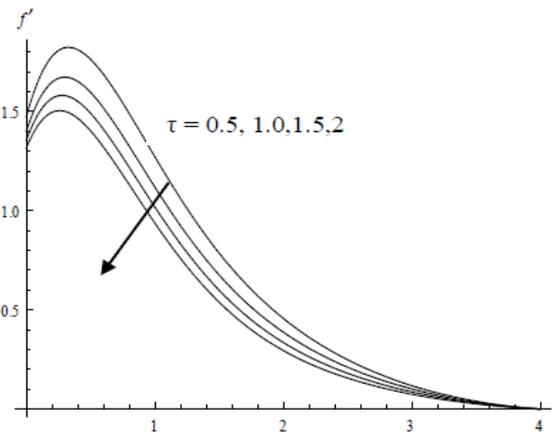


Figure-7a: Variation of  $f'$  with  $\tau$

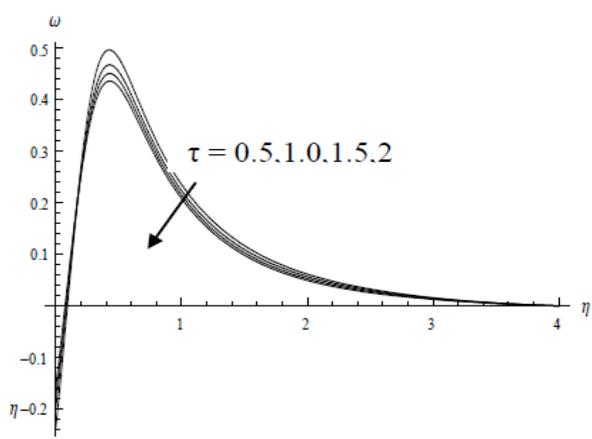


Figure-7b: Variation of  $\omega$  with  $\tau$

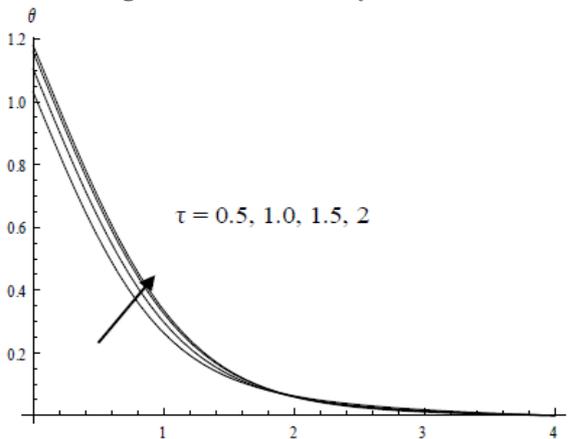


Figure-7c: Variation of  $\theta$  with  $\tau$

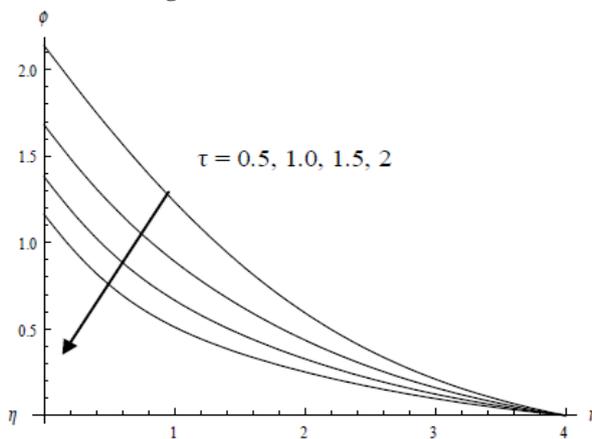


Figure-7d: Variation of  $\phi$  with  $\tau$

**Table-3**

The local Skin-friction (Cf), couple stress (Cw), Nusselt number (Nu), Sherwood number (Sh) for different values of parameters

G1	B1	$\lambda$	A1	B11	Cf	Cw	Nu	Sh
0.2	0.2	0.1	0.1	0.1	2.60812	2.43688	0.96016	0.44462
0.4	0.2	0.1	0.1	0.1	2.42186	2.26287	0.97017	0.46738
0.6	0.2	0.1	0.1	0.1	2.41495	1.70187	0.97005	0.46779
0.8	0.2	0.1	0.1	0.1	2.41291	1.25982	0.96975	0.46821
0.2	0.4	0.1	0.1	0.1	2.65137	2.36947	0.97064	0.46728
0.2	0.6	0.1	0.1	0.1	2.72154	2.33368	0.96201	0.46732
0.2	0.8	0.1	0.1	0.1	2.89404	2.20645	0.95939	0.46737
0.2	0.2	0.3	0.1	0.1	2.34273	0.64294	0.97062	0.46771
0.2	0.2	0.5	0.1	0.1	2.32775	0.28926	0.97009	0.46792
0.2	0.2	0.7	0.1	0.1	2.31964	0.15369	0.96912	0.46802
0.2	0.2	0.1	0.3	0.1	2.48136	2.43469	1.04467	0.46349
0.2	0.2	0.1	-0.1	0.1	2.47821	2.40981	0.99816	0.46576
0.2	0.2	0.1	-0.3	0.1	2.50608	2.45126	1.02711	0.46426
0.2	0.2	0.1	0.1	0.3	2.46645	2.37855	0.9619	0.46772
0.2	0.2	0.2	0.1	-0.1	2.46575	2.40149	1.0066	0.46537
2	0.2	0.2	0.1	-0.3	2.48136	2.43469	1.04467	0.46349

**Table-4**

The local Skin-friction (Cf), couple stress (Cw), Nusselt number(Nu), Sherwood number(Sh) for different values of parameters

Sr/Du	Ec	Rd	Cf	Cw	Nu	Sh
0.6/0.1	0.01	0.5	2.57459	2.3752	1.00258	0.41944
1.0/0.06	0.01	0.5	2.60812	2.43688	0.96016	0.44462
1.5/0.04	0.01	0.5	2.6867	2.51777	0.89209	0.47532
2.0/0.03	0.01	0.5	2.91524	2.68375	0.79142	0.50812
1.0/0.06	0.03	0.5	2.61476	2.40956	0.93423	0.46953
1.0/0.06	0.05	0.5	2.65911	2.49736	0.91902	0.47032
1.0/0.06	0.07	0.5	2.74774	2.60834	0.88415	0.47246
1.0/0.06	0.01	1.5	3.51427	2.97518	0.61455	0.51444
1.0/0.06	0.01	3.5	4.78569	3.72434	0.44829	0.57619
1.0/0.06	0.01	5.0	5.44464	4.15775	0.40066	0.60583

## 6. CONCLUSIONS

The combined effects of suction and thermophoresis on the mixed convective heat and mass transfer of a fluid through a porous medium in the presence of Thermo-Diffusion and Diffusion-Thermo effects were investigated. Similarity transformations are used to transform the resulting partial differential equations into ordinary differential equations and numerically solved by using the Finite element method. The following conclusions have been made from the present problem.

- 1) Higher the radiative heat flux larger the velocity, angular velocity, temperature, Skin friction, Sherwood number and smaller the concentration and Nusselt Number.
- 2) Higher the dissipation larger the velocity, angular velocity, temperature and smaller the Concentration. The Skin friction and Sherwood number increases and Nusselt number reduces with increases Ec.
- 3) Increasing soret parameter Sr(or Decreasing Du) larger the Velocity, angular velocity, concentration, Nusselt number and smaller the Skin friction and Sherwood number.
- 4) The velocity, angular velocity, temperature, Skin friction, Sherwood number enhances , the Concentration, Nusselt number reduces, with increasing A1, B11>0,while a reversed effect is noticed with A1,B11<0.
- 5) An increase in thermophoresis parameter  $\tau$  reduces the velocity, angular velocity, concentration and enhances the temperature.

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